

# PHY-100C FINAL EXAM SOLUTIONS:

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## PHYSICS 100C Final Exam, Thursday, June 7, 8AM-11AM

1. A loop of radius  $R$  has "glued" linear charge density distribution as a function of azimuthal angle  $\varphi$ ,  $\rho(\varphi) = \rho_0 \cos^2 \varphi$ . The loop is spinning with angular frequency  $\omega$ . Find retarded potentials (Lorentz gauge)  $A$  and  $V$  at the center of the loop.

$$V(0, t) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}', t_R)}{R} d\vec{r}' = \frac{1}{4\pi\epsilon_0} \int_0^{2\pi} \frac{\rho_0 \cos^2 \varphi \cdot R \cdot d\varphi}{R}$$

$$= \frac{\rho_0}{4\pi\epsilon_0} \int_0^{2\pi} \cos^2 \varphi \cdot d\varphi = \frac{\rho_0}{4\pi\epsilon_0} \cdot \pi = \frac{\rho_0}{4\epsilon_0}$$

Here  $\int_0^{2\pi} \cos^2 \varphi \cdot d\varphi = \int_0^{2\pi} \frac{1 + \cos 2\varphi}{2} = \frac{\varphi}{2} \Big|_0^{2\pi} + \frac{\sin 2\varphi}{4} \Big|_0^{2\pi} = \pi$

$$\vec{A}(0, t) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}', t_R)}{R} d\vec{r}'$$

But  $|\vec{J} \cdot d\vec{r}'| = \rho \cdot \omega R$  and  $\vec{J}(\varphi) = -\vec{J}(\pi + \varphi)$   
therefore opposite sections cancel and  $\vec{A} = 0$

2. Show that the x-rays (for which frequency much greater than resonance frequencies typically in UV range  $\omega \gg \omega_j$ ) have an index of refraction less than 1, and estimate the angle of total internal reflection for 10 keV x-rays incident on vacuum/metal interface at a grazing incidence. You may assume  $N = 10^{30}$  "free" electrons per  $m^3$  in a metal.

Eq. 9.170:  $n = 1 + \frac{N_m e^2}{2m\epsilon_0} \sum \frac{f_j (\omega_j^2 - \omega^2)}{(\omega_j^2 - \omega^2)^2 + \gamma_j^2 \omega^2}$

for  $\omega \gg \omega_j$ ;

$$\frac{\omega_j^2 - \omega^2}{(\omega_j^2 - \omega^2)^2 + \gamma_j^2 \omega^2} \approx \frac{-\omega^2}{\omega^4 + \gamma_j^2 \omega^2} \approx -\frac{1}{\omega^2}$$

and since  $N_m = \frac{\text{molecules}}{\text{volume}}$ ,  $f_j = \frac{\text{electrons}}{\text{molecule}}$

$$N = N_m \cdot \sum f_j \left( \frac{\text{electrons}}{\text{volume}} \right)$$

$$n = 1 - \frac{Ne^2}{2m\epsilon_0 \omega^2} = 1 - \delta$$

where  $\delta = \frac{Ne^2}{2m\epsilon_0 \omega^2} = 6.8 \cdot 10^{-6}$

(we use  $\omega = \frac{E}{\hbar} = \frac{10 \text{ keV}}{\hbar} = \frac{10^4 \cdot e}{\hbar}$  and  $e^2$  cancels)

vacuum  $n=1$   $\xrightarrow{\text{x-ray } \alpha}$   
 $n=1-\delta$

SNELL'S LAW:  
 $1 \cdot \sin\left(\frac{\pi}{2} - \alpha\right) = (1 - \delta) \cdot \sin \frac{\pi}{2}$   
OR  
 $\cos \alpha \approx 1 - \frac{\alpha^2}{2} = 1 - \delta$

and  $\alpha = \sqrt{2\delta} = 3.7 \cdot 10^{-3}$   
 OR 3.7 mrad OR 2.1 deg

3. a) A loop of radius R (yes, yet another spinning charged loop problem!) is centered at (0,0,0) and is oriented in x-y plane as shown on the Fig. 1. Upper half of the loop has "glued" uniform linear charge density + $\rho$ , and bottom half of the loop has linear charge density - $\rho$ . Find total power radiated if the loop is spinning around its axis (with respect to z axis) at angular frequency  $\omega$ .

- b) How would the answer change if the loop is instead is spun around x axis (at the same frequency)? How about if it was spinning around y-axis?

Dipole moment (along  $\hat{y}$  at  $t=0$  due to symmetry)

$$P = 2 \int_0^\pi \rho \cdot R \cdot \sin\varphi \cdot R \cdot d\varphi = 2\rho R^2 (-\cos\varphi) \Big|_0^\pi = 4\rho R^2$$

A) DURING ROTATION  $\ddot{\vec{p}} = -\omega^2 \vec{p}$

and therefore due to LARMOR FORMULA

$$\text{Power} = \frac{\mu_0 \cdot |\ddot{\vec{p}}|^2}{6\pi c} = \frac{\mu_0 \omega^4 (4\rho R^2)^2}{6\pi c} = \frac{8\mu_0 \omega^4 \rho^2 R^4}{3\pi c}$$

B) SPINNING around x gives the same ANSWER AS A) since  $\vec{p}$  rotates at  $\omega$  with respect to axis  $\perp \vec{p}$ .

SPIN AROUND y-axis has  $\ddot{\vec{p}} = 0$  therefore no dipole radiation.

1. a) Electric and magnetic fields at some region of space are given by two vectors,  $\mathbf{E}$  and  $\mathbf{B}$ , such as  $\mathbf{E} \perp \mathbf{B}$ . Show that there always exists a frame of reference in which  $\mathbf{B}'=0$ , and find its velocity (direction and magnitude), as well as the resulting electric field  $\mathbf{E}'$  in this frame of reference.  
 b) Show that for two arbitrary (no longer mutually perpendicular) vectors  $\mathbf{E}$  and  $\mathbf{B}$ , their scalar product  $\mathbf{E} \cdot \mathbf{B}$  is relativistically-invariant. Comment on existence of frame of reference found in a) when  $\mathbf{E}$  and  $\mathbf{B}$  are not perpendicular.

A) LETS DEFINE  $\hat{y} \parallel \vec{B}$ ,  $\hat{z} \parallel \vec{E}$   
 $\hat{x} \perp \vec{E}, \vec{B}$ , so that  $B_z = 0, E_y = 0$   
 and  $E_x = B_x = 0, B_y = |\vec{B}|, E_z = |\vec{E}|$

In FRAME OF REFERENCE moving along  $\hat{x}$ :

$$E'_x = B'_x = 0$$

$$E'_y = 0$$

$$B'_y = \gamma \left( B_y + \frac{v}{c^2} E_z \right)$$

$$E'_z = \gamma(E_z + vB_y) \quad B'_z = 0$$

(Eq. 12.108)

FOR  $\vec{B} = 0$  (OR  $B_y = 0$ )

$$B_y + \frac{v}{c^2} E_z = 0 \Rightarrow v = -\frac{B_y \cdot c^2}{E_z}$$

$$\text{OR } v = -\frac{|\mathbf{B}| \cdot c^2}{|\mathbf{E}|}$$

$$\begin{aligned} \text{then } |\mathbf{E}| = E'_z &= \gamma(E_z + vB_y) = \gamma\left(|\mathbf{E}| - \frac{|\mathbf{B}|^2 \cdot c^2}{|\mathbf{E}|}\right) \\ &= \gamma |\mathbf{E}| \cdot \left(1 - \frac{|\mathbf{B}|^2 \cdot c^2}{|\mathbf{E}|^2}\right) \quad \text{and} \end{aligned}$$

$$\gamma = \left[1 - \left(\frac{v}{c}\right)^2\right]^{-1/2} = \left(1 - \frac{|\mathbf{B}|^2 \cdot c^2}{|\mathbf{E}|^2}\right)^{-1/2}, \quad \text{so that}$$

$$|\mathbf{E}| = E'_z = |\mathbf{E}| \cdot \left(1 - \frac{|\mathbf{B}|^2 \cdot c^2}{|\mathbf{E}|^2}\right)^{1/2}$$

$$b) \vec{E}' \cdot \vec{B}' = E'_x \cdot B'_x + E'_y \cdot B'_y + E'_z \cdot B'_z$$

using Eq. 12.108:

$$\begin{aligned} \vec{E}' \cdot \vec{B}' &= E_x B_x + \gamma^2 (E_y - vB_z)(B_y + \frac{v}{c^2} E_z) + \\ &+ \gamma^2 (E_z + vB_y) \cdot (B_z - \frac{v}{c^2} E_y) = E_x B_x + \\ &+ \gamma^2 \left( \underbrace{E_y B_y}_{\text{red}} - \cancel{v B_y B_z} + \underbrace{\frac{v}{c^2} E_y E_z}_{\text{green}} - \frac{v^2}{c^2} B_z E_z + \right. \\ &+ \left. \underbrace{E_z B_z}_{\text{green}} + \cancel{v B_y B_z} - \frac{v}{c^2} \cancel{E_y E_z} - \frac{v^2}{c^2} \underbrace{E_y B_y}_{\text{red}} \right) = \\ &= E_x B_x + \underbrace{\gamma^2 \cdot \left(1 - \frac{v^2}{c^2}\right)}_1 \cdot (E_y B_y + E_z B_z) = \\ &= E_x B_x + E_y B_y + E_z B_z = \vec{E} \cdot \vec{B} \end{aligned}$$

If  $\vec{E}$  is not  $\perp \vec{B}$   $\vec{E} \cdot \vec{B} \neq 0$   
 AND  $\vec{E}' \cdot \vec{B}' \neq 0$  so that  $B' \neq 0$   
 IN ALL FRAMES

5. Large Hadron Collider is set to collide two beams of protons travelling at the same speed in opposite directions with energy of 7 TeV per proton (in the laboratory frame). What is the energy of on-coming protons in the frame of reference of one of the beams?

$$m_0 c^2 = 1 \text{ GeV} \quad E = \gamma m_0 c^2 = 7 \text{ TeV}$$

$$p_x = -\gamma m_0 v \quad p_0 = \frac{E}{c}$$

$$p'_0 = \frac{E'}{c} = \gamma(p_0 - \beta \cdot p_x) = \gamma\left(\frac{E}{c} + \beta \gamma m_0 v\right)$$

$$E' = \gamma E + \gamma^2 m_0 v^2 = \gamma\left(\underbrace{E + \gamma m_0 c^2}_E\right) + \gamma^2 m_0 \underbrace{\left(\frac{v^2}{c^2} - 1\right)}_{-1/\gamma^2} \cdot c^2$$

Since  $\gamma = \frac{E}{m_0 c^2}$

$$E' = 2\gamma E - m_0 c^2 = \frac{2E^2}{m_0 c^2} - m_0 c^2$$

(Also:  $E' = (2\gamma^2 - 1) \cdot m_0 c^2$  or  $E' = \frac{2\gamma^2 - 1}{\gamma} \cdot E$ )

In ANY case, since  $\gamma = 7 \cdot 10^3$ ,

$$E' = 1.4 \cdot 10^4 E \quad \text{OR} \quad 9.8 \cdot 10^{16} \text{ eV} \quad \text{OR} \quad 98 \text{ PeV}$$

Peta  $\rightarrow$