

PHYS 100C

Wednesday, March 31, 2010

Lecture #2:

EM in vacuum: (cont'd)

Maxwell equations:

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0} = 0 \quad \text{Gauss}$$

$$\vec{\nabla} \cdot \vec{B} = 0 \quad \text{No MAGN. MONOPOLES}$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad \text{FARADAY}$$

$$\vec{\nabla} \times \vec{B} = \cancel{\mu_0 \vec{J}} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t} \quad \text{AmPere (+MAXWELL)}$$

Coupled 1st order

CURL IT AWAY:

$$\begin{aligned} \nabla \times (\nabla \times \vec{E}) &= \nabla (\nabla \cdot \vec{E}) - \nabla^2 \vec{E} = -\nabla^2 \vec{E} = \\ &= \nabla \times \left(-\frac{\partial \vec{B}}{\partial t} \right) = -\frac{\partial}{\partial t} (\nabla \times \vec{B}) = -\mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2} \end{aligned}$$

$$\begin{aligned} \nabla \times (\nabla \times \mathbf{B}) &= \nabla (\nabla \cdot \mathbf{B}) - \nabla^2 \mathbf{B} = -\nabla^2 \mathbf{B} = \\ &= \nabla \times \left(\mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \right) = \mu_0 \epsilon_0 \frac{\partial}{\partial t} (\nabla \times \mathbf{E}) = \\ &= -\mu_0 \epsilon_0 \frac{\partial^2 \mathbf{B}}{\partial t^2} \end{aligned}$$

Therefore:

$$\begin{aligned} \nabla^2 \mathbf{E} &= \mu_0 \epsilon_0 \frac{\partial^2 \mathbf{E}}{\partial t^2} \\ \nabla^2 \mathbf{B} &= \mu_0 \epsilon_0 \frac{\partial^2 \mathbf{B}}{\partial t^2} \end{aligned}$$

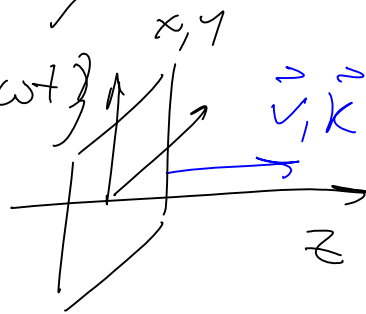
3D wave equations!
with velocity $c^2 = \frac{1}{\mu_0 \epsilon_0}$ SPEED OF LIGHT IN VACUUM $3 \cdot 10^8 \text{ m/s}$

Monochromatic Plane Waves

Waves traveling along z ,
have no x - y dependence:

$$\vec{E}(z,t) = \text{Re} \left\{ \vec{E}_0 \cdot e^{i(kz - \omega t)} \right\}$$

$$\vec{B}(z,t) = \text{Re} \left\{ \vec{B}_0 \cdot e^{i(kz - \omega t)} \right\}$$



Gauss Law:

$$\vec{\nabla} \cdot \vec{E} = 0$$

$$\frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = 0$$

(because no x - y dependence)

$$\frac{\partial E_z}{\partial z} = ik E_0^z = 0 \Rightarrow E_0^z = 0$$

Similarly $B_0^z = 0$

EM waves are transverse!

\vec{E} & \vec{B} vectors \perp propagation

Faraday Law:

$$\vec{\nabla} \times \vec{E} = - \frac{\partial \vec{B}}{\partial t}$$

$$\begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ E_x & E_y & 0 \end{vmatrix} = - \frac{\partial \vec{B}}{\partial t}$$



Along \hat{x} :

$$- \hat{x} \frac{\partial}{\partial z} E_y = - \hat{x} \cdot \frac{\partial B_x}{\partial t}$$

$$ik E_0^y = -i\omega B_0^x \Rightarrow B_0^x = -\frac{k}{\omega} E_0^y$$

Along \hat{y} :

$$\hat{y} \frac{\partial}{\partial z} E_x = - \hat{y} \cdot \frac{\partial B_y}{\partial t}$$

$$ik E_0^x = i\omega B_0^y \Rightarrow B_0^y = \frac{k}{\omega} E_0^x$$

If we pick \hat{x} so that its

along B_0 ; i.e. $B_0^y = 0$

then $E_0^x = 0$ as well

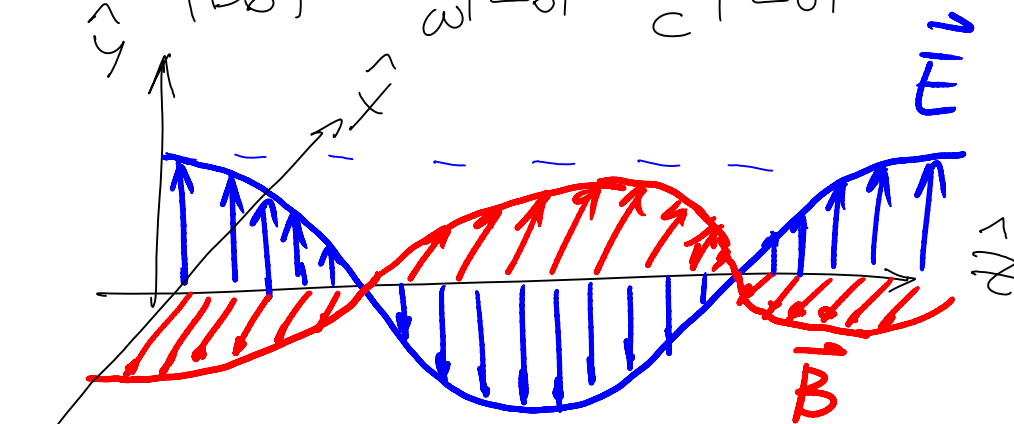
and $\vec{E}_0 \parallel \hat{y}$, $\vec{B}_0 \parallel \hat{x}$,

$$\vec{E}_0 \perp \vec{B}_0$$

OR: $\vec{B}_0 = \frac{k}{\omega} (\hat{z} \times \vec{E}_0)$

\vec{E}_0 and \vec{B}_0 in phase,
and perpendicular to each
other

$$|\vec{B}_0| = \frac{k}{\omega} |\vec{E}_0| = \frac{1}{c} |\vec{E}_0|$$



\vec{E} is in xz plane
 \vec{B} is in yz plane

* What about \hat{z} component?

$$\hat{z} \frac{\partial}{\partial x} E_y - \hat{z} \frac{\partial}{\partial y} E_x = - \frac{\partial B_z}{\partial t} \hat{z}$$

$$B_z = 0 ; \quad \frac{\partial E_y}{\partial x} = \frac{\partial E_x}{\partial y} = 0$$

because waves are PLANAR

(in)

(no x - y dependence)

* Also: AMPERE - Maxwell Eq.
does not produce any new info
Solution above already satisfies.