

# PHYS 100C, LECTURE #15

Tuesday, May 04, 2010

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## \* Jefimenko Equations

↑  
"the other" Oleg

given  $\rho, \mathbf{j}$  for all times, everywhere  
find  $E$  and  $B$ :

$$V(r, t) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(r', t_R)}{r} d\tau'$$

$$A(r, t) = \frac{\mu_0}{4\pi} \int \frac{\mathbf{j}(r', t_R)}{r} d\tau'$$

$$E = -\nabla V - \frac{\partial A}{\partial t} \quad \text{and}$$

$$B = \nabla \times A$$

$$\nabla V = \frac{1}{4\pi\epsilon_0} \int \left( -\frac{\dot{\rho}}{c} \frac{\hat{r}}{r} - \rho \frac{\hat{r}}{r^2} \right) d\tau'$$

(see Lecture #7 notes)

$$\frac{\partial A}{\partial t} = \frac{\mu_0}{4\pi} \int \frac{\dot{\mathbf{j}}}{r} d\tau'$$

$$E = \frac{1}{4\pi\epsilon_0} \int \left( \frac{\rho}{r^2} \hat{r} + \frac{\dot{\rho}}{c r} \hat{r} - \frac{\dot{\mathbf{j}}}{c^2 r} \right) d\tau'$$

Where  $\rho, \dot{\rho}$  and  $\dot{\mathbf{j}}$  evaluated at  $t_R$

Note that only the first term  
is Coulomb Eq. in retarded formalism,  
the terms with  $\dot{\rho}$  and  $\dot{\mathbf{j}}$  are

"UNexpected"! (From naive approach)

Similarly,  $B = \nabla \times A$

$$\nabla \times A = \frac{\mu_0}{4\pi} \int \left[ \frac{1}{r^2} (\nabla \times J) - \dot{y} \times \nabla \left( \frac{1}{r} \right) \right] d\tau'$$

(See  $\nabla \times (f\vec{A})$  in "vector identities")

$\nabla \times J$  has terms like

$$\frac{\partial J_z}{\partial y} = \frac{\partial J_z}{\partial t_R} \cdot \frac{\partial t_R}{\partial y}$$

$$\text{but } t_R = t - r/c \Rightarrow \partial t_R = -\frac{1}{c} \partial r$$

$$\frac{\partial J_z}{\partial y} = \dot{y}_z \cdot \left(-\frac{1}{c}\right) \cdot \frac{\partial r}{\partial y}$$

these are now components of

$$\dot{y} \times (\nabla r) \quad (\text{TRUST BUT VERIFY!})$$

$$\nabla r = \hat{r} \quad \text{and} \quad \nabla \left( \frac{1}{r} \right) = -\frac{\hat{r}}{r^2}$$

$$B = \frac{\mu_0}{4\pi} \int \left( \frac{y}{r^2} + \frac{\dot{y}}{cr} \right) \times \hat{r} \cdot d\tau'$$

Hand-waving (naïve) argument  
(also the best kind of argument!)  
would give us the first term  
for  $B$  in "retarded" time form,  
but not the  $\dot{y}$  term.

Bottom Line: Time retarded  
 ARGUMENT works for potentials  
 $V$  and  $A$ , but not for  $E$  &  $B$ .  
 Expressions for  $E$  &  $B$  (derived  
 from retarded potentials) include  
 "non-obvious" contributions  $\dot{m}_j$  and  $\dot{j}_j$

## \* POINT CHARGES, Liénard-Wiechert POTENTIALS

As before, 
$$V(r, t) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(r', t_r)}{r} d\tau'$$

If we consider a single "point"  
 charge  $q$  moving along some known  
 trajectory  $\vec{w}(t)$ , then a naive  
 person may decide that

$$\rho(r, t_r) = \delta^3(r - w(t_r)) \cdot q$$

and therefore

$$V(r, t) = \frac{1}{4\pi\epsilon_0} \cdot \frac{q}{r}$$

Where  $r$  is retarded distance  
 from  $\vec{w}(t_r)$  to observer.

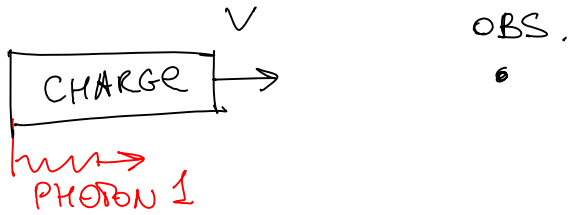
This is WRONG!

(bad day for naive people!)

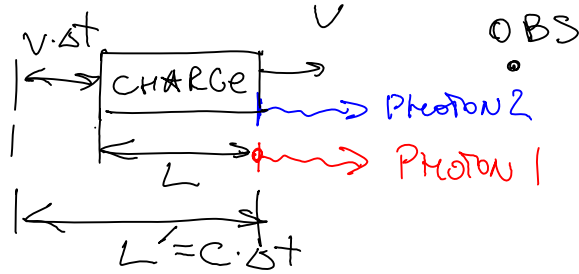
The charge is going to be larger  
 than  $q$ , by a factor  $(1 - \vec{v} \cdot \hat{r}/c)^{-1}$

Why? Geometry/Kinematics:

$t=0$   
CARTOON:



$t = \Delta t$



Observer will "see" PHOTONS 1 & 2 at the same time and decide that actual "length" of charge is

$$L' = L + v \cdot \Delta t = c \cdot \Delta t$$

and since  $\Delta t = \frac{L}{c-v}$

$$L' = L + \frac{Lv}{c-v} = \frac{Lc}{c-v}$$

IF CHARGE IS moving away from observer,

$$L' = \frac{Lc}{c+v} \quad (\text{"shrinkage"})$$

Generally

$$L' = L \cdot \frac{c}{c - v \cdot \cos \theta} = L \cdot \frac{1}{1 - \vec{v} \cdot \hat{z} / c}$$

The "volume" of charge and total

The "volume" of charge and total charge will therefore be multiplied by  $\frac{1}{1 - \vec{v} \cdot \hat{r}/c}$

\* Extreme case:  $v=c$  and moving towards observer.

Observer won't see fields / photons, until all of a sudden they all come in at once! (together with charge)

General Expression for  $V$ :

$$V(r, t) = \frac{1}{4\pi\epsilon_0} \cdot \frac{qc}{rc - \vec{r} \cdot \vec{v}}$$

AND FOR  $A$ , similarly, since  $\vec{j} = \rho \vec{v}$ :

$$A(r, t) = \frac{\mu_0}{4\pi} \cdot \frac{qc\vec{v}}{rc - \vec{r} \cdot \vec{v}} = \frac{\vec{v}}{c^2} \bar{V}(r, t)$$

These are Liénard-Wiechert potentials.